

Fourier Analysis

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Review.

• Thm (Fourier inversion formula)

Let $f \in \mathcal{M}(\mathbb{R})$. Suppose $\hat{f} \in \mathcal{M}(\mathbb{R})$.

Then

$$f(x) = \int_{\mathbb{R}} \hat{f}(\xi) e^{2\pi i \xi x} d\xi$$

• Thm (Plancherel formula).

Let $f \in \mathcal{M}(\mathbb{R})$. Suppose that $\hat{f} \in \mathcal{M}(\mathbb{R})$.

Then

$$\int_{-\infty}^{\infty} |f(x)|^2 dx = \int_{-\infty}^{\infty} |\hat{f}(\xi)|^2 d\xi.$$

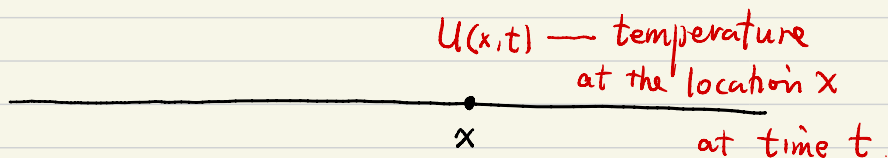
• Schwartz space $\mathcal{S}(\mathbb{R})$.

• $f \in \mathcal{S}(\mathbb{R})$ if $f \in C^{\infty}(\mathbb{R})$ and

$$\sup_{x \in \mathbb{R}} |x^n f^{(l)}(x)| < \infty \text{ for all } n, l \geq 0.$$

• $f \in \mathcal{S}(\mathbb{R}) \Leftrightarrow \hat{f} \in \mathcal{S}(\mathbb{R})$.

§ 5.2 Application (I): Heat equation on \mathbb{R} .



$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, & x \in \mathbb{R}, t > 0 & \textcircled{1} \\ u(x, 0) = f(x), & x \in \mathbb{R} & \textcircled{2} \end{cases}$$

(initial condition)

Let us use the method of Fourier transform to derive a solution for the above heat equation.

Taking Fourier transform on the both sides of $\textcircled{1}$ (in variable x),

we obtain

$$\frac{\partial}{\partial t} \widehat{u}\left(\frac{\xi}{2}, t\right) = (2\pi i \frac{\xi}{2})^2 \widehat{u}\left(\frac{\xi}{2}, t\right).$$

Justification:

$$\text{Set } \hat{u}(\xi, t) = \int_{\mathbb{R}} u(x, t) e^{-2\pi i \xi x} dx$$

Then

$$\begin{aligned} \int_{\mathbb{R}} \frac{\partial u}{\partial t}(x, t) e^{-2\pi i \xi x} dx \\ &= \frac{\partial}{\partial t} \int_{\mathbb{R}} u(x, t) e^{-2\pi i \xi x} dx \\ &= \frac{\partial}{\partial t} \hat{u}(\xi, t). \end{aligned}$$

$$\begin{aligned} \int_{\mathbb{R}} \frac{\partial^2 u(x, t)}{\partial x^2} \cdot e^{-2\pi i \xi x} dx \\ &= (2\pi i \xi)^2 \hat{u}(\xi, t). \end{aligned}$$

Hence

$$\frac{d \hat{u}(\frac{x}{3}, t)}{dt} = -4\pi^2 \frac{x}{3}^2 \hat{u}(\frac{x}{3}, t)$$

So

$$\frac{d \hat{u}(\frac{x}{3}, t)}{dt} + 4\pi^2 \frac{x}{3}^2 \hat{u}(\frac{x}{3}, t) = 0.$$

$$e^{4\pi^2 \frac{x}{3}^2 t} \left(\frac{d \hat{u}(\frac{x}{3}, t)}{dt} + 4\pi^2 \frac{x}{3}^2 \hat{u}(\frac{x}{3}, t) \right) = 0$$

So

$$\frac{d}{dt} \left(e^{4\pi^2 \frac{x}{3}^2 t} \hat{u}(\frac{x}{3}, t) \right) = 0.$$

$$e^{4\pi^2 \frac{x}{3}^2 t} \hat{u}(\frac{x}{3}, t) = A(\frac{x}{3}).$$

$$\Rightarrow \hat{u}(\frac{x}{3}, t) = A(\frac{x}{3}) \cdot e^{-4\pi^2 \frac{x}{3}^2 t}.$$

Taking $t=0$ in the above equation gives

$$\hat{u}(\xi, 0) = A(\xi)$$

Recall that $u(x, 0) = f(x)$. It follows that

$$\hat{u}(\xi, 0) = \hat{f}(\xi).$$

Hence $A(\xi) = \hat{f}(\xi)$. Therefore

$$\hat{u}(\xi, t) = \hat{f}(\xi) \cdot e^{-4\pi^2 \xi^2 t}.$$

Notice that $e^{-4\pi^2 \xi^2 t}$ is the Fourier transform of

$$\mathcal{H}_t(x) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}, \quad x \in \mathbb{R}, t > 0.$$

We call $\{\mathcal{H}_t\}_{t>0}$ the heat kernel on \mathbb{R} as $t \rightarrow 0$.

So

$$\begin{aligned}\widehat{u}(\xi, t) &= \widehat{f}(\xi) \cdot \widehat{H}_t(\xi) \\ &= \widehat{f * H_t}(\xi)\end{aligned}$$

By the Fourier inversion formula, we obtain

$$u(x, t) = f * H_t(x).$$

Next we give a theoretic check.

Thm. Let $f \in S(\mathbb{R})$. Let $u(x, t) = f * H_t(x)$.
Then the following holds:

① $u \in C^\infty(\mathbb{R} \times \mathbb{R}_+)$, where $\mathbb{R}_+ := (0, \infty)$,

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}.$$

② $u(x, t) \rightrightarrows f(x)$ as $t \rightarrow 0$

③ $\int_{\mathbb{R}} |u(x, t) - f(x)|^2 dx \rightarrow 0$ as $t \rightarrow 0$

Pf: ① Since $f, \mathcal{H}_t \in S(\mathbb{R})$,
 So are \widehat{f} and $\widehat{\mathcal{H}_t}$. It follows that
 $\widehat{f} \cdot \widehat{\mathcal{H}_t} \in S(\mathbb{R})$. That is, $\widehat{f * \mathcal{H}_t} \in S(\mathbb{R})$.

Hence

$$f * \mathcal{H}_t \in S(\mathbb{R}).$$

So by the inversion formula,

$$\begin{aligned} u(x,t) &= \int_{\mathbb{R}} \widehat{f * \mathcal{H}_t}(\xi) \cdot e^{2\pi i \xi x} d\xi \\ &= \int_{\mathbb{R}} \widehat{f}(\xi) \cdot e^{-4\pi^2 \xi^2 t} e^{2\pi i \xi x} d\xi \end{aligned}$$

In the following, we show $\frac{\partial u}{\partial t}$ exists

$$\begin{aligned} \frac{\partial u}{\partial t}(x,t) &= \lim_{\Delta t \rightarrow 0} \frac{u(x, t+\Delta t) - u(x, t)}{\Delta t} \\ &= \lim_{\Delta t \rightarrow 0} \int_{\mathbb{R}} \widehat{f}(\xi) \frac{e^{-4\pi^2 \xi^2 (t+\Delta t)} - e^{-4\pi^2 \xi^2 t}}{\Delta t} \cdot e^{2\pi i \xi x} d\xi \end{aligned}$$

$$= \lim_{\Delta t \rightarrow 0} \int_{\mathbb{R}} \hat{f}(\xi) \cdot e^{-4\pi^2 \xi^2 t} \cdot \frac{e^{-4\pi^2 \xi^2 \Delta t} - 1}{\Delta t} \cdot e^{2\pi i \xi x} d\xi$$

Notice that

$$\left| \frac{e^{-4\pi^2 \xi^2 \Delta t} - 1}{\Delta t} \right| \leq \text{const} \cdot |\xi|^2 \quad (\text{using the mean value theorem})$$

So

$$\left| \hat{f}(\xi) e^{-4\pi^2 \xi^2 t} \cdot \frac{e^{-4\pi^2 \xi^2 \Delta t} - 1}{\Delta t} \cdot e^{2\pi i \xi x} \right| \leq \left| \hat{f}(\xi) \right| \cdot |\xi|^2 \quad (\text{Notice that } \hat{f}(\xi) \cdot \xi^2 \in \mathcal{S}(\mathbb{R}))$$

By the dominated convergence thm,

$$\frac{\partial u(x,t)}{\partial t} = \int_{\mathbb{R}} \hat{f}(\xi) \cdot e^{-4\pi^2 \xi^2 t} (-4\pi^2 \xi^2) \cdot e^{2\pi i \xi x} d\xi$$

In a similar way, we can show $u \in C^\infty(\mathbb{R} \times \mathbb{R}_+)$.

Moreover,

$$\begin{aligned}\frac{\partial^2 u}{\partial x^2} &= \int_{\mathbb{R}} \hat{f}(\xi) \cdot e^{-4\pi^2 \xi^2 t} \frac{\partial^2}{\partial x^2} \left(e^{2\pi i \xi x} \right) d\xi \\ &= \int_{\mathbb{R}} \hat{f}(\xi) e^{-4\pi^2 \xi^2 t} \cdot (-4\pi^2 \xi^2) \cdot e^{2\pi i \xi x} d\xi \\ &= \frac{\partial u}{\partial t}\end{aligned}$$

This proves ①.